
1. Maxwell's insight

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8 points

Maxwell's laws are as follows:

- electrostatica:

$$- \oint_{\text{kring}} \vec{E} \cdot d\vec{l} = 0$$

$$- \Phi_E \equiv \oint_{\text{oppervlak}} \vec{E} \cdot d\vec{\sigma} = \frac{Q^{\text{omsloten}}}{\epsilon_0}$$

- magnetostatica:

$$- \oint_{\text{kring}} \vec{B} \cdot d\vec{l} = \mu_0 I^{\text{omsloten}}$$

$$- \Phi_B \equiv \oint_{\text{oppervlak}} \vec{B} \cdot d\vec{\sigma} = 0$$

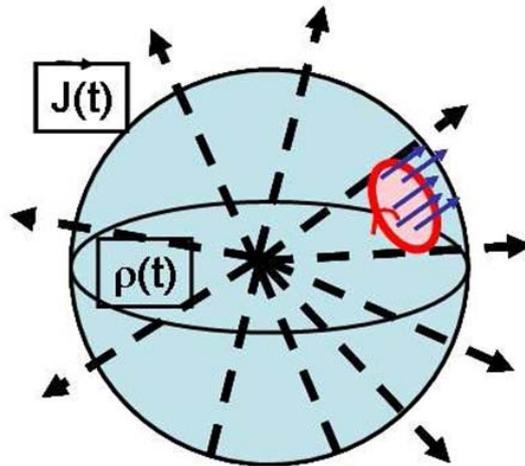


Figure 1.1: The amount of charge in the origin

In the origin of our coordinate system, we find an amount of electrical charge: $Q(t) = Q_0 e^{-\frac{t}{\tau}}$ where t is the time, τ is a constant and Q_0 is the charge at time $t = 0$.

1. Calculate, using Gauss' law, the electrical field that is induced by the charge in the origin. Indicate with a drawing which 'Gaussian surface' you've used.

The electrical charge in the origin can't just disappear. We presume that the charge moves away (flows) in a radial direction.

2. Give an expression for the electric current $I(t)$ and the density current \vec{J} . (Be aware that the density current is a vector and therefore, it has a direction.)

3. Argument, with use of a symmetry argument, there can't be a magnetic field parallel to the surface of the imaginary ball in figure 1.1.

4. Argument, if necessary with the Ampère's circuital law, that there's a magnetic field parallel to the surface of the ball. You could be inspired by the Ampère loop which is included in figure 1.1.

The answers given in **3.** and **4.** are, as you might have noticed, contradictory. Luckily, we all know that Maxwell solved this contradiction in a brilliant fashion, by adding a term to Ampère's circuital law.

5. Give Ampère's circuital law as an integral, with Maxwell's addition. Show that, with this extra term, the magneticfield at **4.** will disappear as well. (Hint: use answers found at **1.** and **2.**)

2. Noise

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9 points

In the northern part of the Dutch province Groningen, extensive research is done to noise caused by F16's. The terrain where the investigation takes place, is open. At a certain moment, the F16 is spotted, flying at constant height towards the observers. Directly after seeing the F16, a recording device is turned which measures the spectrum of sound during a period of time. The frequency of the loudest tone is shown in the graph, as a function of time. The minimum height at which an F16 is allowed to fly, is known to be 1000 ft. We neglect the absorption of sound.

The following is known:

The speed of sound through air is 343ms^{-1}

$1\text{ft} = 0.3048\text{m}$

Determine, using the data in graph 2.1

1. The speed of the F16,
2. The frequency of the loudest tone in the recorded spectrum,
3. Whether or not the F16 flies above 1000 ft,
4. The horizontal distance at which the F16 was first noticed.

Please clarify how you calculate your results.

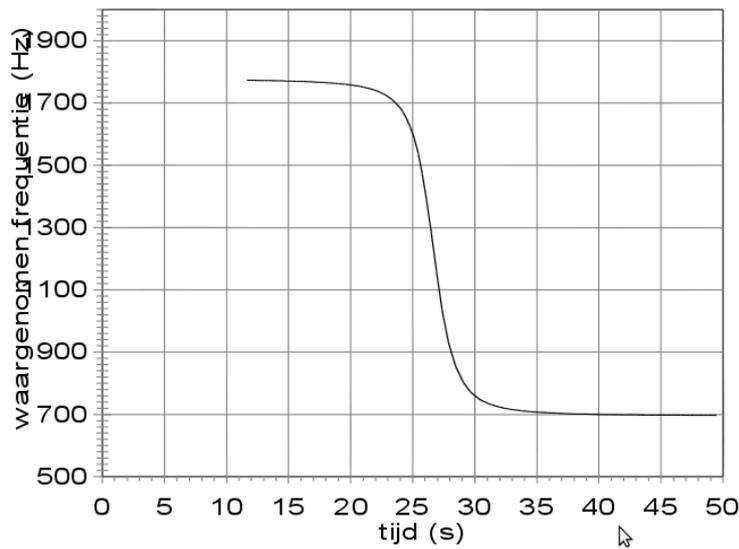


Figure 2.1: The recorded frequency (Hz) as a function of time (s)

3. Thermoelectric transport and heating

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11 punten

Consider an electrical conductor, which has a temperature difference ΔT and a voltage difference ΔV .

1. Start with a temperature difference equal to zero ($\Delta T = 0$) and presume that the electrical current I is determined by Ohm's Law: $I = \Delta V R^{-1}$, with R the resistance of the conductor. Use the first law of Thermodynamics, $dU = dQ + \mu dN = TdS + \mu dN$ (where dU : change of energy, dQ the heat, μ : the chemical potential, dN : the change of particles T the temperature and dS the change of entropy), to show that the total heat, produced by the electric current per time unit, is given by:

$$\frac{dQ}{dt} = \frac{\Delta V^2}{R} \quad (3.1)$$

2. Let us now discuss the situation $\Delta V = 0$, $\Delta T = 0$. Now we find an electric current I , as well as an heat wave I_Q through the conductor. Presume that:

$$I = \frac{\Delta V}{R} - \frac{S\Delta T}{R} \quad (3.2)$$

and

$$I_Q = \frac{P\Delta V}{R} - \kappa' \Delta T \quad (3.3)$$

Here S is the so called Seebeck coefficient and P the Peltier coefficient. Furthermore κ' is the thermal conductivity at $\Delta V = 0$. Use the First Law to show that:

$$\frac{dQ}{dt} = \frac{\kappa' \Delta T^2}{T} + \frac{\Delta V^2}{R} - \frac{P\Delta V\Delta T}{RT} - \frac{S\Delta V\Delta T}{R} \quad (3.4)$$

3. Use the Second Law of Thermodynamics to show that the transport coefficients must satisfy: $\kappa' R \geq PS$

4. A rotating cylinder

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10 points

In 1908 Paul Ehrenfest proposed a paradoxal case to physicists who were working on the, then new, special theory of relativity. The problem gave birth to lots of discussion and inspired Albert Einstein to further develop his theory.

Ehrenfest had the following in mind:

A rigid body in the shape of a cylinder, height H , radius R and circumference O , ($O = 2\pi R$), is initially at rest and is then put into a circular movement around its axis. From a certain point the cylinder rotates with a constant angular speed. An observer who moves along with the cylinder will measure a radius R' and a perimeter O' .

1. What do you expect, according to the theory of special relativity, to be the ratio between R' and O' .
2. What do you think is problematic in this particular case?
3. Try to give your own analysis on what's happening to the cylinder.

5. Flow in a pipe

Han Velthuis, TNO Fluid Dynamics

8 points

Consider a long, round pipe with water flowing through. The inner radius of the pipe is $R = 0.025m$. The length of the pipe is $L = 1km$. We presume that the flow through the pipe is stationary, i.e. that the velocity profile of the water in the pipe isn't changing as a function of time nor as a function of the coordinate in the flow direction. The velocity profile $u(r)$ is purely a function of the radius r , where the radius is measured from the axis of the pipe. The variable u depicts the velocity of water. The pressure drop over the pipe length L is represented by variable Δp .

- Beneath a critical value the flow will be laminar. Solving the theoretical equations of flow (Navier-Stokes, 1822) will give a parabole as velocity profile. The speed of water will be zero near the side ($r = R$), and maximal at the axis ($r = 0$).

- Above a critical value the flow will be turbulent. The turbulent velocity profile is often described with the so called 'power-law', where the velocity profile is proportional with $(Y/R)^c$. Here y is the distance to the side and c is a constant (say $c = \frac{1}{7}$).

- The critical value is given by the so called Reynolds number $Re = \frac{\rho u_{ave} D}{\mu}$, named after the famous British scientist Osborne Reynolds (1842-1912). Here ρ is the density of water, D is the inner diameter of the pipe and μ is the dynamic viscosity of water.

The average speed of water by the tube cross section is given by u_{ave} . The transition from a laminar flow to a turbulent flow usually takes place at a critical Reynolds number of approximately $Re \approx 2300$.

Water often behaves like a Newtonian fluid. When that's the case, the viscose shear stress τ , which is exercised by the water on the side of the pipe, is given by the dynamic viscosity of water multiplied by the velocity gradient of the water at the side of the pipe, $\tau = \mu \frac{du}{dr}|_{r=R}$.

Questions

1. Derive an expression for the loss of pressure Δp along the pipe as a function of the shear stress τ at the side of the wall.

2-a. Construct a graph of this pressure drop along the pipe as a function of the average velocity of water, u_{ave} , with the reach of velocity $u_{ave} \in [0, 0.1]ms^{-1}$. Use a BINAS (a Dutch book with important properties of elements and basic compositions) to find the correct properties of water.

(Hint: In order to determine the shear stress, one has to determine the shape of the velocity profile as a function of the radius r).

2-b. Why is it impossible to construct the entire graph?

3-a. An empirical formula of turbulent pressure drop along the pipe Δp is given by:

$$\Delta p = \frac{0.3164}{(Re)^{0.25}} \frac{L}{2R} \frac{1}{2} \rho u_{ave}^2 \quad (5.1)$$

What's the graph now?

3-b. What's striking about this graph?

6. Fluorescence from a Single Molecule

G.A. Blab, Universiteit Utrecht
10 punten

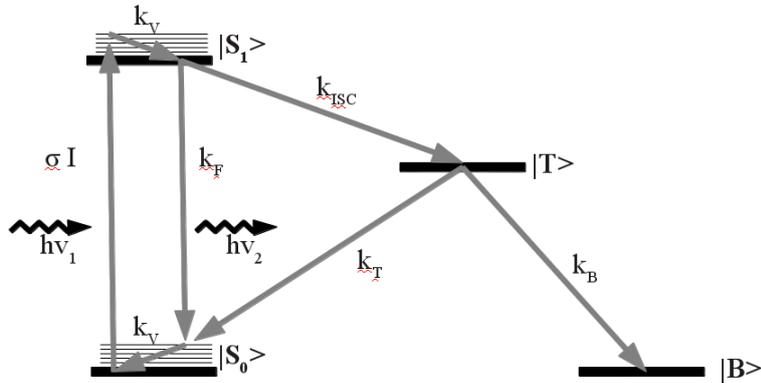


Figure 6.1: Simplified Jablonski diagram of a fluorescent molecule. A photon $h\nu_1$ can excite the molecule from its ground state $|S_0\rangle$ to an excited state $|S_1\rangle$. From there, the molecule can emit a photon with lower energy $h\nu_2$ and go back to the ground state, or transition through a triplet (dark) state $|T\rangle$ back to the ground state, or bleach to become a photo product $|B\rangle$. When asked for a numerical answer, use the following rate constants: radiative $k_F = 3 \times 10^8 \text{ s}^{-1}$; non-radiative $k_V = 1 \times 10^{12} \text{ s}^{-1}$, $k_{ISC} = 1 \times 10^7 \text{ s}^{-1}$, $k_T = 2 \times 10^5 \text{ s}^{-1}$, $k_{bl} = 1 \times 10^4 \text{ s}^{-1}$.

1 Find an analytical expression for the measured fluorescence life time τ_F , that is the life time of the excited state $|S_1\rangle$ after absorption of a photon $h\nu_1$. From the general expression, make the assumption $k_V \gg k_F, k_{ISC}$. How does this effective fluorescence life time τ_F differ from $(k_F)^{-1}$?

2 Ignoring for the moment the irreversible bleaching, what is the expected time constant for one excitation emission cycle τ_{cycle} ? What does this mean for the number of photons a molecule can emit per unit time? You may use the approximation $k_{ISC} \gg k_T$ in your calculations.

3 Find the excitation intensity (expressed in $W\text{cm}^{-2}$) needed for a photon emission(!) rate of $5 \times 10^4 \text{ s}^{-1}$ from a single molecule. To find the absorption cross section σ , you may assume that the conditions are such that the Beer-Lambert law is valid and that the molar extinction coefficient of the fluorophore is $\epsilon = 80\,000 \text{ M}^{-1} \text{ cm}^{-1}$ at the excitation wavelength $\lambda_{\text{exc}} = 488 \text{ nm}$.
